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Crystal structure, spectroscopy and thermodynamic properties of $M^{\rm l} {\rm VWO}_6(M^{\rm l}$ – Li, Na)

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1. Introduction

Compounds with brannerite structure [\[15\]](#page-8-0) containing alkali metal atoms can be used as materials for independent supply source. There is information about LiVMoO $_6$ and LiVWO $_6$ and their use in lithium-ion batteries [\[1\].](#page-8-0) Besides materials on basis of $\rm M^1$ VWO₆ compounds can be used as catalysts due to content of transition metals.

The object of our investigation is complex oxides containing vanadium, tungsten and alkali metals. In order to understand the relations among structure and energetics in a fundamental sense, and to correlate reactivity at high temperature and materials compatibility in applications, thermodynamic data are essential. Despite the extensive interest in vanadium tungsten oxides, there is no any thermodynamic data and other physicochemical data are presented in bare number of publications. Therefore, detailed investigations of these compounds are especially important.

The goals of this work include detailed structural and spectroscopic investigations, studies of thermal stability, calorimetric determination of the temperature dependence of the heat capacity $C_p^{\circ} = f(T)$ of LiVWO₆ and NaVWO₆ from 7 to 640 K, detection of the possible phase transitions, and calculation of the standard thermo-

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ABSTRACT

In the present work lithium (sodium) vanadium tungsten oxides with brannerite structure is refined by the Rietveld method (space group $C2/m$, $Z = 2$). IR and Raman spectroscopy was used to assign vibrational bands and determine structural particularities. The diffuse reflectance spectra allow to calculate bandgap for M^1 VWO₆(M^1 – Li, Na). The temperature dependences of heat capacity have been measured first in the range from 7 to 350 K for these compounds and then between 330 and 640 K, respectively, by precision adiabatic vacuum and dynamic calorimetry. The experimental data were used to calculate standard thermodynamic functions, namely the heat capacity $C_p^{\circ}(T)$, enthalpy $H^{\circ}(T)$ – $H^{\circ}(0)$, entropy $S^{\circ}(T) - S^{\circ}(0)$ and Gibbs function $G^{\circ}(T) - H^{\circ}(0)$, for the range from $T \to 0$ to 640 K. The differential scanning calorimetry was applied to measure decomposition temperature of compounds under study. $©$ 2009 Elsevier Inc. All rights reserved.

> dynamic functions $C_p^{\circ}(T)$, $H^{\circ}(T)$ - $H^{\circ}(0)$, $S^{\circ}(T)$ - $S^{\circ}(0)$ and $G^{\circ}(T)$ - $H^{\circ}(0)$ in the range from $T\rightarrow 0$ to 640 K.

2. Experimental

2.1. Sample

Lithium (sodium) vanadium tungsten oxide samples were prepared by the solid-state reaction between tungsten oxide, vanadium oxide and alkali metal nitrate, respectively [\[1\]](#page-8-0). The synthesis was performed in a porcelain crucible, into which the reaction mixture with the atomic ratio1Li (Na):1W:1V was loaded. The mixture was calcined at 873 K for 20 h, undergoing regrinding every 5 h. For structural investigations, X-ray diffraction patterns of LiVWO₆ and NaVWO₆ samples were recorded on a Shimadzu X-ray diffractometer XRD-6000 (CuK α radiation, geometry θ -2 θ) in the 2θ range from 10° to 120° with scan increment of 0.02°. Rietveld analysis and structure refinement [\[2\]](#page-8-0) were carried out using RIETAN-94 software [\[3\]](#page-8-0). The X-ray data and estimated impurity content (0.5–1 wt%) in the substances led us to conclude that the studied sample were individual crystalline compounds.

2.2. Apparatus and measurement procedure

To measure the heat capacity C_p ^o of the tested substances in the range from 7 to 350 K a BKT-3.0 automatic precision adiabatic

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vacuum calorimeter with discrete heating was used. The calorimeter design and the operation procedure were described earlier [\[4,5\]](#page-8-0). The calorimeter was tested by measuring the heat capacity of high-purity copper and reference samples of synthetic corundum and K-2 benzoic acid. The analysis of the results showed that measurement error of the heat capacity of the substance at helium temperature was within $\pm 2\%$, then it decreased to \pm 0.5% as the temperature was rising to 40 K, and was equal to \pm 0.2% at T > 40 K. Temperatures of phase transitions can be determined with the error of $+0.02$ K.

To measure the heat capacity of the sample between 330 and 650 K an automatic thermoanalytical complex (ADKTTM) – a dynamic calorimeter operating by the principle of triple thermal bridge was used [\[6,7\].](#page-8-0) The device design and the measurement procedure of the heat capacity, temperatures and enthalpies of phase transitions were demonstrated in detail in the abovementioned papers. The reliability of the calorimeter operation was checked by measuring the heat capacity of the standard sample of synthetic corundum as well as the thermodynamic characteristics of fusion of indium, tin and lead. As a result, it was found that the calorimeter and the measurement technique allow one to obtain the heat capacity values of the substances in solid and liquid states with the maximum error of \pm 1.5% and the phase transition temperatures within ca. \pm 0.5 K. Since the heat capacity of the examined compound was also measured between 330 and 350 K in the adiabatic vacuum calorimeter with the error of \pm 0.2% and the conditions of measurements in the dynamic device were chosen so that in the above temperature interval the C_p ^o values measured with the use of both calorimeters coincided, it was assumed that at $T>395$ K the heat capacity was determined with the error of 0.5–1.5%. The data on the heat capacity of the objects under study were obtained in the range from 330 to 650 K at the average rate of heating of the calorimeter and the substances of 0.0333 K/s.

Thermal behavior was carried out with DSC Labsys from Setaram in a platinum crucible ranging from 293 to 1173 K (heating rate 0.167 K/s).

Polycrystalline infrared spectra were measured with a Biorad 575C FT-IR spectrometer in KBr suspension for the 1200– 400 cm $^{-1}$ region and in Nujol suspension for the 500–50 cm $^{-1}$ region. FT-Raman spectra were measured using BRUKER 110/S spectrometer. Excitation was performed with a 1064 nm line of a $YAG:Nd^{3+}$ laser. Both IR and Raman spectra were recorded with a spectral resolution of 2 cm $^{-1}$.

Room-temperature diffuse reflectance spectra were obtained in the 200–2500 nm spectral region using a Cary-Varian 5E UV– vis–near-IR spectrophotometer and diffuse reflectance accessory.

3. Results and discussion

3.1. Crystal structure

The structure of LiVWO₆ and NaVWO₆ was refined assuming space group $C2/m$. The initial model included the atomic coordinates in the structure of NaVMoO $₆$ [\[1,8\]](#page-8-0). The details of the</sub> X-ray diffraction experiment and structure refinement data are listed in Table 1. [Fig. 1](#page-2-0) represents the measured, simulated and difference X-ray diffraction patterns for LiVWO $_6$ and NaVWO $_6$, as well as a pattern of lines corresponding to reflection maxima. There is a good agreement between the measured and simulated patterns. [Table 2](#page-2-0) lists the coordinates of the atoms and their isotropic thermal parameters. The refined model yielded positive isotropic thermal parameters B for all atoms. In the case of $NaVWO₆$, we calculated thermal parameters for all atoms types, whereas for lithium counterpart we equated their values for all

Table 1

Details of the X-ray diffraction experiment and the results of the structure refinement for M^1 VWO₆ (M^1 – Li, Na).

Definition of reliability factors R_{WD} and R_{p} are given as follows: R_{WD} = $\{(\Sigma w_i[y_{i\text{obs}}-y_{i\text{calc}}]^2/(\Sigma w_i[y_{i\text{obs}}]^2)\}^{1/2}; R_p = (\Sigma |y_{i\text{obs}}-y_{i\text{calc}}|)/(\Sigma y_{i\text{obs}})$

light atoms in order to get positive isotropic thermal parameters in defined model. This procedure is used in powder X-ray diffraction quite often. [Table 3](#page-2-0) lists the interatomic distances and valence angles in the structure and [Fig. 2](#page-3-0) represents a fragment of the $Li(Na)VWO₆$ structure.

As we can see from the obtained X-ray diffraction data the compounds under investigation have quasi-layered structure. The VWO₆ layers consist of distorted V(W)-octahedra jointed by edges. The alkali metal atoms are located in the interlayer space forming $M^{I}O_{6}$ -polyhedra. The $M^{I}O_{6}$ -polyhedra are composed of three oxygen atoms from one layer and three similar oxygen atoms from neighboring layer. In case of NaVWO $₆$, the coordina-</sub> tion of sodium ions is close to octahedral, although there is a slight compression of the octahedron along axial axis in this structure and its tendency towards tetragonal–pyramidal environment. In LiVWO₆, the LiO₆-polyhedron is tetragonal bipyramid, in which the difference between equatorial and axial bond distances is equal to 0.472 Å . The alkali metal coordination polyhedra are joined by edges forming linear (one-dimensional) chains.

Analysis of the unit cell parameters shows that the lattice parameters (a and b) of Li(Na)VWO₆ are very similar and these values slightly increase when moving from LiVWO $_6$ to NaVWO $_6$ by 0.77% and 0.18%, respectively. In contrast to this behavior, the unit cell along crystallographic direction c expands significantly (by 9.2%) when Li⁺ is replaced by Na⁺. This behavior is connected with smaller compression of Na-coordination polyhedron than Li-polyhedron leading to increase of the unit cell volume by 10.5%.

In order to establish the crystal-chemical borders in alkali metal vanadium tungsten oxides, we have analyzed several compounds with brannerite structure. It was determined that only negligible quantity of compounds could crystallize in this structure type. $M^1VA^{VI}O_6$ (M^1 – Li, Na, K; A^{VI} – Mo, W) [\[1,8,17\],](#page-8-0) $M^{II}V_2O_6$ (M^{II} – Mg, Ca, Zn, Cd, Mn) [\[9–13\]](#page-8-0) and $M^{IV}Ti_2O_6$ (M^{IV} – U, Th) [\[15,14\]](#page-8-0) compounds are among them. Consequently, there are small atoms (0.54 $A \le r \le 0.61$ Å) with +4, +5, +6 oxidation numbers in the layers and average size atoms $(0.72 \text{ Å} \leq r \leq 1.38 \text{ Å})$ with +1, +2, +4 oxidation numbers in the interlayer space. These criteria allow to suppose that complex vanadium tungsten oxides with large alkali metals could not crystallize in brannerite structure. Such conclusion was confirmed in paper [\[16\]](#page-8-0), where authors indicated $CsVWO₆$ as compound with pyrochlore structure and cubic system (space group F d3m).

Fig. 1. Fragments of (1) observed, (2) simulated, and (3) difference X-ray diffraction patterns for LiVWO₆ and NaVWO₆, (4) Bragg reflections. The simulated pattern is shifted relative to the observed pattern.

Table 2

Coordinates and isotropic thermal parameters of atoms in the structure of M^1 VWO $_6$ $(M^I - Li, Na)$.

Atom	Site	χ	y	z	Occ	$B(\AA^2)$
LiVWO ₆						
Li	2a	Ω	Ω	Ω		0.6(1)
V	4i	0.8159(2)	$\bf{0}$	0.3543(2)	0.5	1.40(6)
W	4i	0.8159(2)	$\bf{0}$	0.3543(2)	0.5	1.40(6)
O(1)	4i	0.982(1)	Ω	0.291(1)		0.6(1)
O(2)	4i	0.690(1)	Ω	0.111(1)		0.6(1)
O(3)	4i	0.293(1)	Ω	0.413(1)		0.6(1)
NaVWO ₆						
Na	2a	Ω	Ω	Ω		1.6(2)
V	4i	0.8180(1)	$\bf{0}$	0.3663(2)	0.5	0.25(3)
W	4i	0.8180(1)	$\bf{0}$	0.3663(2)	0.5	0.25(3)
O(1)	4i	0.9752(8)	Ω	0.310(1)		1.0(2)
O(2)	4i	0.6711(9)	Ω	0.135(1)		0.9(2)
O(3)	4i	0.3094(8)	$\mathbf{0}$	0.438(1)		0.3(2)

3.2. Raman and IR

The structure of LiVWO₆ and NaVWO₆ is composed of $MeO₆$ ($Me = W, V$) and LiO₆ (NaO₆) polyhedra [\[18\]](#page-8-0). Factor group analysis predicts for this structure $(C2/m \, (C_{2h}^3)$ space group) $8A_g + 4B_g +$ $5A_u+10B_u$ Brillouin zone center modes. These modes can be subdivided into A_u+2B_u and $2A_g+B_g+A_u+2B_u$ translations of the Li⁺ $(Na⁺)$ and Me⁶⁺ cations, respectively. It should be remembered, however, that among the translational modes two modes (A_u+2B_u)

 $a - atoms$ sites O: $1 - x$, 0, $1 - z$.

Fig. 2. Fragment of the structure of $Li(Na)VWO₆$.

belong to the acoustic branches. The remaining modes $(6A_{\sigma}+3B_{\sigma}+$ $3A_u+6B_u$) correspond to vibrations of the oxygen atoms. The A_g and B_g modes are Raman active, and the A_u and B_u modes are IR active.

The Raman spectra show presence of very strong bands at 968 and 965 cm $^{-1}$ for LiVWO₆ and NaVWO₆, respectively (see Fig. 3). The $968\,\mathrm{cm^{-1}}$ band of LiVWO $_6$ was assigned previously to vibrations of the terminal $V=O$ groups [\[19\]](#page-8-0). Since tungsten and vanadium ions occupy statistically the same sites in the crystal structure, we assign this band to vibrations of the terminal Me =0 groups. The corresponding IR bands are observed as clear doublets near $956-973$ cm^{-1} (see Table 4 and Fig. 3). The IR spectra show also presence of weak bands near 936–934 $\rm cm^{-1}$ which can be most likely assigned to some overtones or combination bands.

The bands of LiVWO₆ in the regions 809–886 and 576– 754 cm^{-1} were assigned previously to asymmetric stretching vibrations of the V–O–V bonds running parallel to the b-axis and $V₂O₂$ groups, respectively [\[19\]](#page-8-0). Closer inspection of the structure shows, however, that all MeO $_6$ octahedra running parallel to the b-axis share edges, not corners. As a result no single oxygen bridges of the type V–O–V are present in the structure but the Me ions are joined along the b-axis through double oxygen bridges MeOOMe with the shorter (Me–O) and longer (Me \cdots O) distances equal to 1.905 and 2.175 Å or 1.903 and 2.152 Å according to Ref. [\[18\]](#page-8-0) and our present X-ray data, respectively. This double oxygen bridge will be denoted as type I. These adjacent chains of the $MeO₆$ octahedra running parallel to the b-axis also share edges forming double oxygen bridges with the Me–O and Me \cdots O distances of 1.687 and 2.379 or 1.760 and 2.401 \AA according to Ref. [\[18\]](#page-8-0) and our present X-ray data, respectively. This bridge will be denoted as type II. We assign the 809–886 and 576–754 $\rm cm^{-1}$ bands of LiVWO₆ to stretching vibrations of asymmetric bridge II and symmetric bridge I, respectively. The corresponding bands for bridge I of NaVWO $₆$ are observed at very similar wavenumbers but</sub> the bands assigned by us to bridge II are significantly shifted towards higher wavenumbers (up to $12\,\mathrm{cm}^{-1}$). Inspection of the structure shows that the oxygen atoms which take part in formation of bridge I are well separated from the Li⁺ (Na⁺) ions whereas the oxygen atoms taking part in formation of bridge II have close contacts with Li⁺ ions. As a result the substitution of Li⁺ by Na+ should have significant influence only on bridge II. Indeed, our X-ray results show that whereas the Me–O and Me \cdots O distances for bridge I are very similar for both LiVWO₆ and NaVWO $_6$ (see [Table 3](#page-2-0)), the Me–O bond in bridge II is much shorter for NaVWO₆ (1.677 Å) than LiVWO₆ (1.760 Å). The observed shift of bands in the 809–886 cm $^{-1}$ region towards higher wavenumbers

Fig. 3. Raman (a), mid-IR (b) and far-IR (c) spectra of LiVWO₆ and NaVWO₆.

Table 4

Raman and IR wavenumbers for LiVWO $_6$ and NaVWO $_6$ together with the proposed assignment.

LiVWO ₆		NaVWO ₆		Assignment
Raman	IR	Raman	IR	
968s	973m	965s	968sh	$v(Me = 0)$
	956m		956m	
	936w		934 _w	Overtone or combination band
861sh	886s	874w	896m	v(MeOMeO) type II
		847sh	870 _w	
823m	809s	832m	823m	
740m	754 _w	737m	749w	v(MeOMeO) type I
719sh				
	625s		628s	
	576s		573m	
		519 _{vw}		Overtone or combination band
446m	448sh	463m	453w	δ(MeOMeO) type II
	383vw		401 _w	δ (MeOMeO) type I and II
	363w		368w	
324w	318s	322w	332s	
279sh	295s	282sh	297s	
266m		264m		
	251m		182m	δ (MeOMeO)+T'(Li ⁺) (or T'(Na ⁺)
238m		231m		δ(MeOMeO) type I
206w		216sh		
	195 _w		159w	$T'(Me^{6+})+T'(Li^+)$ (or $T'(Na^+)$)
		183m		$T'(Me^{6+})$
146m		144m		
118w		110w		

for NaVWO $_6$ supports, therefore, our assignment of these bands to vibrations of bridge II and it is consistent with significant shortening of the Me–O bond. Amdouni et al. [\[19\]](#page-8-0) also suggested that the weak Raman band near 519 cm $^{-1}$ can be assigned to symmetric stretching of the Me–O–Me bridges. It is very well known that symmetric vibrations give usually rise to strong Raman bands and weak IR bands whereas asymmetric vibrations are usually observed as strong IR and weak Raman bands. Very weak intensity of the Raman band at 519 cm $^{-1}$ and its relatively large bandwidth suggest that this band should be assigned to an overtone or a combination band, rather than to the symmetric stretching mode.

Amdouni et al. [\[19\]](#page-8-0) also assigned the band of LiVWO₆ at 446 cm^{-1} to symmetric stretching of the Me–O–Me bridges. Numerous studies of crystals containing $WO₆$ or $VO₆$ octahedra showed that stretching modes are usually observed above $500 \, \text{cm}^{-1}$ but in some cases can also be observed in the 400– 500 cm^{-1} range [20-22]. We propose to assign the band at $446\,\mathrm{cm^{-1}}$ to bending mode of the double oxygen bridge II because this band is well separated from the stretching modes and similarly as in the case of bridge II stretching modes, this band is significantly shifted towards higher wavenumbers for $NaVWO₆$. The IR and Raman bands in the 295–401 and 206–324 $\rm cm^{-1}$ range, respectively, can be assigned to bending modes of the both oxygen bridges. Amdouni et al. also suggested that translations of Li⁺ ions contribute significantly to the IR bands at 252 and 295 cm^{-1} . Replacement of $Li⁺$ (atomic mass 6.94) by much heavier Na⁺ ions (atomic mass 22.99) should lead to wavenumber decrease of the pure translational mode by a factor of $(22.99/6.94)^{0.5} = 1.82$. Our results show, however, the 295 $\rm cm^{-1}$ band of LiVWO₆ shifts to 297 cm $^{-1}$ for NaVWO₆. Therefore, the contribution of the alkali metal translational motions to this band is negligible. In contrast to this behavior, the IR band at 251 $\rm cm^{-1}$ for LiVWO₆ shifts to 182 cm^{-1} for NaVWO₆ indicating that it has significant contribution of alkali metal translation. The bands below $200\,\mathrm{cm}^{-1}$ were assigned previously to the lattice modes [\[19\]](#page-8-0). We confirm this conclusion and assign the Raman bands at 183, 144–146 and 110–118 cm^{-1} to three expected translations of the Me⁶⁺ ions.

3.3. Absorption

The diffuse reflectance spectra of the LiVWO₆ and NaVWO₆ samples are shown in Fig. 4. The spectra exhibit strong increase of absorption at wavelengths lower than 560 and 510 nm for LiVWO₆ and NaVWO $_6$, respectively, due to the bandgap transition. For indirect bandgap material, the absorption coefficient satisfies the equation $(\alpha h v)^{1/2} = A(hv - E_g)$, where α , v , A and E_g are absorption coefficient, light frequency, proportionality coefficient and bandgap energy, respectively [\[23\]](#page-8-0). The bandgap obtained by extrapolation of the plots (insets in Fig. 4) of $(\alpha h v)^{1/2}$ versus $h v$ are 1.73 eV for LiVWO₆ and 2.24 eV for NaVWO₆, respectively. These results indicate that the both materials have a bandgap suitable for photocatalitic applications in the visible range. The observed shift of the optical bandgap from 1.73 eV for LiVWO₆ to 2.24 eV for NaVWO $₆$ is consistent with the change of the sample color from</sub> light brown to nearly yellow. The bandgap values are significantly lower than those obtained for WO_3 (2.6–3.5 eV) [\[24,25\].](#page-9-0) However, the former studies of $(\mathsf{WO}_3)_{1-x}(\mathsf{V}_2\mathsf{O}_5)_x$ system indicated that the bandgap decreases with increasing concentration of V_2O_5 [\[24\]](#page-9-0). It is also worth noting that for LiVWO $_6$ a broad band is observed at about 1570 nm. The origin of this band is not clear.

3.4. Heat capacity

The C_p ^o measurements were carried out between 7 and 640 K. The masses of the sample loaded in the calorimetric ampoules of

Fig. 4. Diffuse reflectance spectra of LiVWO $_6$ (a) and NaVWO $_6$ (b).

the BKT-3.0 and ADKTTM devices were, respectively, 0.8662 and 1.7536 g for LiVWO₆ and 1.8008 and 2.0893 g for NaVWO₆. In the BKT-3.0 calorimeter, 154 (for LiVWO₆) 187 (for NaVWO₆) experimental C_p ^o values were obtained in two series of experiments. The heat capacity of the sample varied from 30% to 70% of the total heat capacity of calorimetric ampoule + substance over the range between 7 and 640 K. The experimental points of C_p° in the temperature interval between $T = (7 \text{ and } 640)$ K were fitted by means of the least-squares method and polynomial equations (Eqs. (1)–(3)) of the C_p° versus temperature have been obtained. The corresponding coefficients (A, B, C, etc.) are given in [Tables 5](#page-5-0) [and 6](#page-5-0)

$$
C_p^0 (T) = A + B \cdot (T/30) + C \cdot (T/30)^2 + D \cdot (T/30)^3
$$

+
$$
E \cdot (T/30)^4 + F \cdot (T/30)^5 + G \cdot (T/30)^6
$$

+
$$
H \cdot (T/30)^7 + I \cdot (T/30)^8 + J \cdot (T/30)^9
$$
 (1)

$$
C_p^0 (T) = A + B \cdot \ln(T/30) + C \cdot {\ln(T/30)}^2 + D \cdot {\ln(T/30)}^3
$$

+
$$
E \cdot {\ln(T/30)}^4 + F \cdot {\ln(T/30)}^5 + G \cdot {\ln(T/30)}^6
$$

+
$$
H \cdot {\ln(T/30)}^7 + I \cdot {\ln(T/30)}^8 + J \cdot {\ln(T/30)}^9
$$
 (2)

Table 5

Coefficients in the fitting polynomials for LiVWO₆.

Table 6

Coefficients in the fitting polynomials for NaVWO₆.

$$
\ln C_p^0 (T) = A + B \cdot \ln(T/30) + C \cdot \{\ln(T/30)\}^2 + D \cdot \{\ln(T/30)\}^3
$$

$$
+ E \cdot \{\ln(T/30)\}^4 + F \cdot \{\ln(T/30)\}^5 + G \cdot \{\ln(T/30)\}^6
$$

$$
+ H \cdot \{\ln(T/30)\}^7 + I \cdot \{\ln(T/30)\}^8 + J \cdot \{\ln(T/30)\}^9 \tag{3}
$$

Their root mean square deviation from the averaging $C_p^{\;\;\;\;\;\;\;\;\;\;\;\;\;\;\;} f(T)$ curve was $\pm 0.15\%$ in the range $T = (6-30)K$, $\pm 0.075\%$ from $T = (25-150)$ K, ± 0.15 % between $T = (130$ and 350) K and ± 0.5 % over the range from $T = (370-650)$ K.

The experimental values of the molar heat capacity of $Li(Na)VWO₆$ over the range from 7 to 640K and the averaging C_p ^o = f(T) plot are presented in Fig. 5 and [Tables 9 and 10](#page-7-0). The heat capacity C_p ^o of this substance gradually increases with rising temperature and does not show any peculiarities until 640 K.

3.5. Standard thermodynamic functions

To calculate the standard thermodynamic functions ([Tables 7](#page-6-0) [and 8](#page-6-0)) of the lithium (sodium) vanadium tungsten oxides, their C_p° values were extrapolated from the starting temperature of the measurement (approximately 7 K) to 0 K by Debye's function of heat capacity:

$$
C_p^o = nD(\theta_D/T) \tag{4}
$$

where D is the symbol of Debye's function, $n = 6$ and $\theta_D(LiVWO_6) = 215.4 \text{ K}$ as well as $\theta_D(NaVWO_6) = 200.4 \text{ K}$ are specially selected parameters. Eq. (4) with the above parameters describes the experimental C_p° values of the compound between 7 and 12 K with the error of \pm 1.69% and \pm 1.77%, respectively.

Fig. 5. Temperature dependences of heat capacity of LiVWO₆ and NaVWO₆.

Table 7

Thermodynamic functions of crystalline LiVWO₆; M(LiVWO₆) = 337.7289 g mol $^{-1}$, $p^{\circ} = 0.1$ MPa.

Table 8

Thermodynamic functions of crystalline NaVWO₆; M(NaVWO₆) = 353.77767 g mol^{-1} , $p^{\circ} = 0.1$ MPa.

 $($ kJ mol $^{-1})$ (kJ mol $^{-1}$)

 $-[G^{\rm o}(T)-H^{\rm o}(0)]$

Table 9 Experimental values of isobaric heat capacities of LiVWO₆, J/(Kmol)
(M(LiVWO₆) = 337.7289 g mol⁻¹).

Table 10

Experimental values of isobaric heat capacities of NaVWO₆, J/(Kmol)
(M(NaVWO₆) = 353.77767 gmol⁻¹).

^a Each third value of C_p^o .

^a Each third value of C_p^o .

In calculating the functions it was assumed that Eq. (4) reproduces the C_p° values of Li(Na)VWO₆ at $T<$ 7K with the same error. The calculations of $H^0(T) - H^0(0)$ and $S^0(T) - S^0(0)$ were made by the numerical integration of $C_p^{\circ} = f(T)$ and $C_p^{\circ} = f(\ln T)$ curves, respectively, and the Gibbs function $G^{\circ}(T) - H^{\circ}(0)$ was estimated from the enthalpies and entropies at the corresponding temperatures [\[26\].](#page-9-0) It was suggested that the error of the function values was \pm 1% at T < 40 K, \pm 0.5% between 40 and 80 K, \pm 0.2% in the range from 80 to 390 K and $+1.5%$ between 390 and 650 K.

The absolute entropies of Li(Na)VWO₆ ([Tables 7 and 8](#page-5-0)) and the corresponding simple substances W(cr), V (cr), Li (cr), Na (cr) [\[27\]](#page-9-0) and $O₂$ (g) [\[28\]](#page-9-0) were used to calculate the standard entropy of formation of the compounds under study at 298.15 K, $\Delta_f S^\circ$ (298.15, LiVWO₆, cr) = -532.1 ± 0.7 J K $^{-1}$ mol $^{-1}$ and $\varDelta_{{\rm f}}$ S°(298.15, NaVWO₆, $cr) = -543.6 \pm 0.7 \,\mathrm{J} \,\mathrm{K}^{-1} \,\mathrm{mol}^{-1}.$

In order to correlate obtained thermodynamic and structural data, it is important to apply to the fractal version of Debye's theory. The values of the fractal dimension $D -$ the most significant parameter of the fractal version of Debye's theory of heat capacity of solids [\[29,30\]](#page-9-0) – were determined from the experimental data on the heat capacity of compounds under investigation. The values of D were estimated using the technique described in works [\[31,32\]](#page-9-0) with the use of Eq. (5):

$$
C_{v} = 3D(D+1)kN\gamma(D+1)\xi(D+1)(T/\theta_{\text{max}})^{D}
$$
\n(5)

where N is the number of atoms in a formula unit, k the Boltzmann constant, $\gamma(D+1)$ the γ -function, $\zeta(D+1)$ the Riemann ξ -function and θ_{max} is the characteristic temperature. According to Ref. [\[33\]](#page-9-0), $D = 1$ corresponds to the solids with chain structure, $D = 2$ to those with a layered structure, and $D = 3$ [\[34\]](#page-9-0) to those with a spatial structure. For the tested compounds $D = 2.4$ for LiVWO₆ and $D = 2.5$ for NaVWO₆ at $T > 30$ K correspond to the intermediate structure between layered and frame structure, that good coordinates with our X-ray diffraction data.

3.6. Differential scanning calorimetry

Joint application of the high temperature X-ray diffraction and thermal analysis (TG-DTA) made it possible to establish some peculiarities of processes taking place in the compounds under investigation during heating. Fig. 6 represents DTA curves of $M^{\rm l} {\rm VWO}_6$ ($M^{\rm l}$ – Li, Na), where we can see endothermic effects at 1022 K for LiVWO₆ and 1031 K for NaVWO₆. These effects are connected with incongruent melting. It is worth noting that melting temperatures of LiVWO₆ and Li₂WO₄ ($T_m = 1013$ K), as well as NaVWO₆ and Na₂WO₄ ($T_m = 1018$ K) are practically equal. The compounds containing sodium and having formulas $M^{I}VO_{3}$, $\mathsf{M}_{2}^{\text{I}}\mathsf{WO}_{4}$ and $\mathsf{M}^{\text{I}}\mathsf{VWO}_{6}$ possess higher melting temperatures than the lithium counterparts by about 5–10 K. This fact is associated with greater M^I-anion sublattice bond energy.

4. Conclusions

Structure of LiVWO₆ and NaVWO₆ is refined by the Rietveld method (space group $C2/m$, $Z = 2$) with R-factor 8.38% and 5.13%,

Fig. 6. Plot of the DTA-signal against temperature for LiVWO₆ and NaVWO₆.

respectively. Also we used IR and Raman spectroscopy for determination of structural particularities. The diffuse reflectance spectra allow to calculate bandgap for $M^{!}$ VWO $_6$ ($M^{!}$ – Li, Na). These results indicate that the both materials have a bandgap suitable for photocatalytic applications in the visible range. The temperature dependences of heat capacity have been measured ranging from 7 to 640 K. The experimental data were used to calculate standard thermodynamic functions for the range from $T\rightarrow 0$ to 640 K. It was shown by the differential scanning calorimetry that these compounds are incongruent melting.

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